

APMA 1941G Homework 10 Solutions
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Problem 1

This problem refers to Example 4, the Earth-Moon Spacecraft Problem. Show that the limiting velocity V_∞ satisfies

$$(V_\infty)^2 = \frac{\tan(\alpha)}{b}.$$

Use the formula relating V_∞ and V_0 , as well as the definition of e , the formula for e in terms of α , and finally the definition of b .

Proof. We defined $e = |r_0||V_0|^2 - 1$. We also have $e = \frac{1}{\cos(\alpha)}$ and $b = \left(\frac{\cos \alpha + 1}{\sin \alpha}\right) |r_0|$. In lecture, we found that

$$(V_\infty)^2 = |V_0|^2 - \frac{2}{|r_0|}.$$

First, we can rewrite

$$|V_0|^2 = \frac{e + 1}{|r_0|} = \frac{\frac{1}{\cos \alpha} + 1}{|r_0|} = \frac{1 + \cos \alpha}{|r_0| \cos \alpha}.$$

Then

$$\begin{aligned} (V_\infty)^2 &= |V_0|^2 - \frac{2}{|r_0|} \\ &= \frac{1 + \cos \alpha}{|r_0| \cos \alpha} - \frac{2 \cos \alpha}{|r_0| \cos \alpha} \\ &= \frac{1 - \cos \alpha}{|r_0| \cos \alpha}. \end{aligned}$$

We can also re-arrange to find that

$$|r_0| = \frac{b}{\left(\frac{\cos \alpha + 1}{\sin \alpha}\right)},$$

so

$$\begin{aligned} (V_\infty)^2 &= \frac{1 - \cos \alpha}{\cos \alpha} \cdot \frac{\left(\frac{\cos \alpha + 1}{\sin \alpha}\right)}{b} \\ &= \frac{\left(\frac{1 - \cos^2 \alpha}{\sin \alpha}\right)}{\cos \alpha b} \\ &= \frac{\left(\frac{\sin^2 \alpha}{\sin \alpha}\right)}{\cos \alpha b} \\ &= \frac{\sin \alpha}{\cos \alpha b} \\ &= \frac{\tan \alpha}{b} \end{aligned}$$

as claimed.

Problem 2: Reynold's Equation for Sliders

Consider a slider that you put in oil, where x is the position between 0 and 1 and the corresponding height is $h(x)$. This slider exerts a certain pressure p^ϵ which satisfies the ODE

$$\begin{aligned} -\epsilon((h^3)p^\epsilon p^\epsilon_x)_x &= (p^\epsilon h)_x \\ p^\epsilon(0) &= p^\epsilon(1) = 1 \end{aligned} \tag{1}$$

Here $p^\epsilon = p^\epsilon(x)$ and $x \in (0, 1)$, and $h = h(x) : [0, 1] \rightarrow (0, \infty)$ is a given height, with $h(1) = 1$ and ϵ is a "viscosity coefficient." We expect there to be a boundary layer near $x = 1$, and our goal is to find a good approximation p^* of p^ϵ .

(a) Outer solution, near $x = 0$

Apply the ansatz $p^\epsilon(x) = p^0(x) + \epsilon p^1(x) + \dots$. Compare the $O(1)$ terms and get an equation for $p^0(x)$. Impose $p^0(0) = 1$, solve for p^0 in terms of $h(0)$ and $h(x)$.

Solution. Substituting the ansatz into (1), we obtain

$$-\epsilon(h^3(p^0 + \epsilon p^1 + \dots)(p_x^0 + \epsilon p_x^1 + \dots))_x = ((p^0 + \epsilon p^1 + \dots)h)_x = 0.$$

We can see that the only $O(1)$ term will be $(p^0 h)_x$ on the right-hand side, which gives

$$(p^0 h)_x = 0.$$

Integrating this in x one time, we have

$$p^0(x)h(x) = C.$$

Then

$$p^0(x) = \frac{C}{h(x)}.$$

Imposing the boundary condition $p^0(0) = 1$, we have that

$$1 = \frac{C}{h(0)} \implies C = h(0).$$

Thus, we can express

$$\boxed{p^0(x) = \frac{h(0)}{h(x)}}.$$

(b) Inner solution, near $x = 1$

Let $y = \frac{x-1}{\epsilon}$ and $\bar{p}_\epsilon(y) = p_\epsilon(x)$ and $\bar{h}(y) = h(x)$. Rewrite (1) in terms of \bar{p}_ϵ and \bar{h} , and apply the ansatz $\bar{p}_\epsilon(y) = \bar{p}_0(y) + \epsilon \bar{p}_1(y) + \dots$, and moreover Taylor expand the function $\bar{h}(y) = h(1 + \epsilon y)$. Compare the $O(1)$ -terms and recall $h(1) = 1$ to obtain

$$\bar{p}_0(\bar{p}_0)_y + \bar{p}_0 = -A.$$

Solve for y in terms of \bar{p}_0 using separation of variables.

Note. At some point, it may be useful to note that

$$\frac{\bar{p}_0}{A + \bar{p}_0} = 1 - \frac{A}{A + \bar{p}_0}.$$

Impose the condition $\bar{p}_0(0) = 1$, which is the same as $p^0(1) = 1$, and ultimately obtain that

$$-y = \bar{p}_0 - A \ln \left| \frac{A + \bar{p}_0}{A + 1} \right| - 1$$

where A is to be determined.

Proof. Rewriting,

$$-\epsilon((h^3)\bar{p}^\epsilon \bar{p}_x^\epsilon)_x = (\bar{p}^\epsilon h)_x$$

Now Taylor expanding \bar{h} , we have that $\bar{h}(y) \approx h(1) + \epsilon y h'(1)$. Since $h(1) = 1$, $\bar{h}(y) \approx 1 + \epsilon y h'(1)$. Considering our ansatz and using the chain rule,

$$\bar{p}_x^\epsilon = \frac{1}{\epsilon} (\bar{p}_y^0 + \epsilon \bar{p}_y^1 + \dots)$$

so we have, after moving the $1/\epsilon$ outside the derivative, transforming the x -derivatives into y -derivatives, and multiplying through by ϵ ,

$$-\epsilon((h^3)(\bar{p}^0 + \epsilon \bar{p}^1)(\bar{p}_y^0 + \epsilon \bar{p}_y^1))_y = (h(\bar{p}^0 + \epsilon \bar{p}^1 + \dots))_y$$

Using the Taylor expansion for h , we see that $O(1)$ terms are:

$$-(\bar{p}^0 \bar{p}_y^0)_y = \bar{p}_y^0$$

Re-arranging and integrating in y , we find that there is a constant A such that

$$\bar{p}^0 \bar{p}_y^0 + \bar{p}^0 = -A.$$

Rewriting this,

$$\frac{d\bar{p}_0}{dy} = \frac{-A - \bar{p}_0}{\bar{p}_0}.$$

Using separation of variables, we have that

$$\begin{aligned} -y &= \int \frac{\bar{p}_0}{A + \bar{p}_0} d\bar{p}_0 \\ &= \int \left(1 - \frac{A}{A + \bar{p}_0} \right) d\bar{p}_0 \\ &= \bar{p}_0 - A \int \frac{1}{A + \bar{p}_0} d\bar{p}_0 \\ &= \bar{p}_0 - A \ln |A + \bar{p}_0| + C \end{aligned}$$

Since $p_0(1) = 1$, $\bar{p}_0(0) = 1$, so $0 = 1 - A \ln |A + 1| + C$. Thus

$$-y = \bar{p}_0 - A \ln \left| \frac{A + \bar{p}_0}{A + 1} \right| - 1, \quad (2)$$

so we have the desired implicit formula for $\bar{p}_0(y)$ in terms of A .

(c) Matching

Here, the matching condition is

$$\lim_{x \rightarrow 1} p^0(x) = \lim_{y \rightarrow -\infty} \bar{p}^0(y).$$

Since we want a finite answer in (2), the only way this works is if the term inside of the logarithm goes to 0^+ , and this is achieved only if \bar{p}_0 goes to $-A$. Use this to solve for A and rewrite (2) in terms of the answer A . Finally, calculate your composite solution $p^*(x)$; leave the \bar{p}_0 term as $\bar{p}_0 \left(\frac{x-1}{\epsilon} \right)$.

Solution. Since

$$\lim_{x \rightarrow 1} p^0(x) = \frac{h(0)}{h(1)} = h(0),$$

$-A = h(0)$. Thus (2) becomes

$$-y = \bar{p}_0 + h(0) \ln \left| \frac{-h(0) + \bar{p}_0}{-h(0) + 1} \right| - 1.$$

The composite solution is given by

$$p^*(x) = p_0(x) + \bar{p}_0(y) - \text{common part}.$$

Based on how we have been defining the common part in class, and the above, the common part is $A = -h(0)$. Thus, leaving the \bar{p}_0 term as $\bar{p}_0 \left(\frac{x-1}{\epsilon} \right)$, we obtain

$$\boxed{p^*(x) = \frac{h(0)}{h(x)} + \bar{p}_0 \left(\frac{x-1}{\epsilon} \right) + h(0).}$$